

Seesaw and effective operators

Florian Bonnet

Laboratoire de Physique Théorique
Université Paris XI

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GDR Neutrino 2007, LAPP Annecy

1 Introduction

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2 Seesaw and effective theory

- Various realization of the Seesaw
- Effective theory

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3 Seesaw Type II

- Model and neutrino mass
- Low-energy effects

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- 3 Seesaw Type II

ν are massless in the SM

- No Dirac mass: $m_D(\bar{\nu}_L\nu_R + \bar{\nu}_R\nu_L)$
 - no right-handed neutrino $\nu_R \rightarrow$ no $Y\tilde{H}_L\bar{\nu}_R + h.c.$

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 - SM renormalizable \rightarrow no dim 5 operator
$$\mathcal{O}_{Weinberg} = Y_{\alpha\beta} \left(\overline{I}_{L\alpha}^c \tilde{H}^* \right) \left(\tilde{H}^\dagger I_{L\beta} \right) + h.c.$$

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\rightarrow Neutrinos are massless

Experimental observations of solar, atmospheric and terrestrial neutrinos showed their oscillations : *Strumia 06*

$$\Delta m_{12}^2 = (8.0 \pm 0.3) \times 10^{-5} \text{eV}^2$$

$$|\Delta m_{23}^2| = (2.5 \pm 0.2) \times 10^{-3} \text{eV}^2$$

$$\tan^2 \theta_{12} = 0.45 \pm 0.05$$

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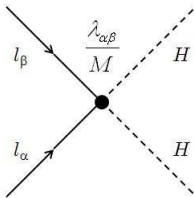
→ ν are massive

flavor mixing in the leptonic sector

Proof that there must be a **Physic Beyond the SM**

Adding heavy fields

The heavy fields manifest in the low energy theory by non-renormalisable operators.



$$\frac{\lambda_{\alpha\beta}}{M} \overline{l_{L\alpha}^c} l_{L\beta} \tilde{H}^* \tilde{H}^\dagger \rightarrow \frac{\lambda_{\alpha\beta}}{M} v^2 \overline{\nu_{L\alpha}^c} \nu_{L\beta}$$

Majorana mass, violation of lepton number

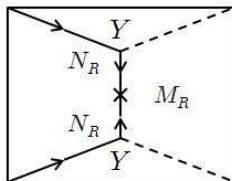
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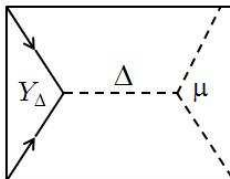
Various realization of the seesaw



Type I

N_R singlet fermion

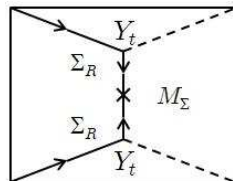
Gell-Man, Ramond,
Slansky
Yanagida
Mohapatra, Senjanovic



Type II

Δ scalar triplet

Georgi, Glashow,
Nussinov
Mohapatra, Senjanovic
Schechter, Valle
Ma, Sarkar



Type III

Σ_R triplet fermion

Ma, Hambye et al.
Bajc, Senjanovic
Abada, Antusch,
Biggio, F.B.
Gavela, in progress...

Development of the Lagrangian into a serie of **non-renormalizable** operators when the heavy fields are integrated out

$$\mathcal{L}_{SM} + \mathcal{L}_{seesaw} \rightarrow \mathcal{L}_{eff} = \mathcal{L}_{SM} + \frac{1}{\Lambda} \mathcal{L}^{d=5} + \frac{1}{\Lambda^2} \mathcal{L}^{d=6} + \dots$$

Development of the Lagrangian into a series of **non-renormalizable** operators when the heavy fields are integrated out

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$$\mathcal{O}_{Weinberg}$$
$$m_\nu$$

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\downarrow \downarrow

$\mathcal{O}_{Weinberg}$ $???$

m_ν

81 operators constructed with the SM fields

(Buchmuller & Wyler, 86)

Exemple : Type I

$$\mathcal{L}_{N_R} = i\overline{N_R}\not{\partial}N_R - \left(\overline{I_L}\tilde{\phi}Y_\nu N_R + \overline{N_R}Y_\nu^\dagger\tilde{\phi}^\dagger I_L\right) - \frac{1}{2}\left(\overline{N_R^c}MN_R + \overline{N_R}M^*N_R^c\right)$$

Effective Lagrangian

$$\mathcal{L}_{\text{eff}} = \mathcal{L}_{SM} + \frac{1}{M}\mathcal{L}^{d=5} + \frac{1}{M^2}\mathcal{L}^{d=6} + \dots$$

$$\mathcal{O}_{\text{Weinberg}} \rightarrow m_\nu \simeq Y_\nu Y_\nu^\dagger \frac{v^2}{M}$$

$$\frac{1}{M^2}\mathcal{L}^{d=6} = \left(\overline{I_L}\tilde{\phi}\right) i\not{\partial} Y_\nu \frac{1}{M^2} Y_\nu^\dagger \left(\tilde{\phi}^\dagger I_L\right) \rightarrow \text{Renormalization of kinetic energy}$$

\Rightarrow **Unitarity violation**

Broncano, Gavela, Jenkins 02

Antusch, Biggio, Fernandez-Martinez, Gavela, Lopez-Pavon 06



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Model of Seesaw de type II

$$\Delta = \begin{pmatrix} \Delta^{++} \\ \Delta^+ \\ \Delta^0 \end{pmatrix} \sim (1, 3, 2), \quad L_\Delta = -2$$

Yukawa term :

$$Y_{\Delta ij} \overline{(l_L)^c}_{ia} (l_L)_{jb} (i\tau_2 \tau_\alpha)_{ab} \Delta^\alpha + h.c$$

Scalar terms:

$$\begin{aligned} & \mu \phi_a^t \phi_b (i\tau_2 \tau_\alpha) (\Delta^\dagger)^\alpha + h.c. \\ & - M_\Delta^2 \Delta^\dagger \Delta - \frac{1}{2} \lambda_2 (\Delta^\dagger \Delta)^2 \\ & - \lambda_3 (\phi^\dagger \phi) (\Delta^\dagger \Delta) \end{aligned}$$

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Masse :

$$m_\nu = Y_\Delta \frac{\mu}{M_\Delta^2} \nu^2$$

2 different scales μ, M_Δ

possibility for $Y_\Delta \sim \mathcal{O}(1)$

$$M_\Delta \sim 1 \text{ TeV} \quad (\mu \sim 100 \text{ eV})$$

d=6 operators

in collaboration with Abada, Gavela, Hambye

$$\frac{1}{2M^2} Y_{\Delta ij} Y_{\Delta kl}^\dagger \left(\bar{l}_{Li} \gamma^\mu l_{Lk} \right) \left(\bar{l}_{Lj} \gamma_\mu l_{Li} \right) \rightarrow \text{Lepton Flavor Violation, constrains}$$

$$\left. \begin{aligned} & -2 \frac{\mu^2}{M_\Delta^4} \partial_\mu (\phi^\dagger \phi) \partial^\mu (\phi^\dagger \phi) \\ & 2\lambda_3 \frac{\mu^2}{M_\Delta^4} (\phi^\dagger \phi)^3 \\ & 4 \frac{\mu^2}{M_\Delta^4} [\phi^\dagger D_\mu \phi]^\dagger [\phi^\dagger D_\mu \phi] \end{aligned} \right\} \rightarrow \text{EW precision data, coupling to bosons}$$

$$-2 \frac{\mu^2}{M_\Delta^4} (\phi^\dagger \phi) \left\{ Y_e \bar{l} e_R \phi + Y_d \bar{q} d \phi - Y_u \bar{q} i \tau_2 u \phi + h.c. \right\} \rightarrow \text{top.....}$$

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The only contribution to ρ is coming from $4 \frac{\mu^2}{M_\Delta^4} [\phi^\dagger D_\mu \phi]^\dagger [\phi^\dagger D_\mu \phi]$

$$\delta\rho = \frac{2v^2\mu^2}{M_\Delta^4} \rightarrow \left| \frac{\mu}{M_\Delta^2} \right| < 3 \times 10^{-5} \text{ GeV}^{-1}$$

to compare with

$$m_\nu = Y_\Delta \frac{\mu}{M_\Delta^2} v^2$$

For $m_\nu \sim 1 \text{ eV}$ we get:

$$\frac{\mu}{M_\Delta^2} : \sim 10^{-15} \text{ GeV}^{-1} \rightarrow \sim 10^{-5} \text{ GeV}^{-1}$$

$$Y_\Delta : 1 \rightarrow 10^{-10}$$

Scalar sector : ILC (1)

$$a_1 \mathcal{O}_{\phi,1} = 4 \frac{\mu^2}{M_\Delta^4} \cdot \frac{1}{2} \partial_\mu (\phi^\dagger \phi) \partial^\mu (\phi^\dagger \phi) ; \quad a_2 \mathcal{O}_{\phi,2} = 6 \lambda_3 \frac{\mu^2}{M_\Delta^4} \cdot \frac{1}{3} (\phi^\dagger \phi)^3$$

Barger & al. 03

- Modification of the scalar potential
- Renormalization of the kinetic term

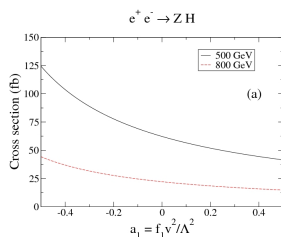
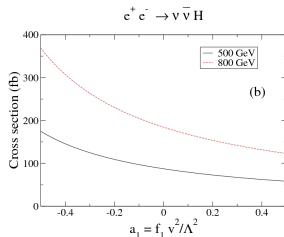
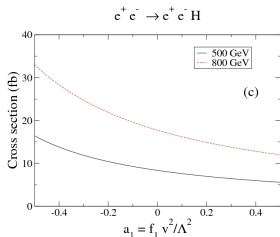
Consequences :

- Modification of HWW, HZZ, HHWW, HHZZ and HHH and HHHH couplings
- Modification of the **single Higgs and Higgs pair production cross sections** at ILC and CLIC

Scalar sector : ILC (2)

$$\text{Single Higgs production} / a_1 = 4 \frac{\mu^2}{M_\Delta^4}$$

Barger & al. 03



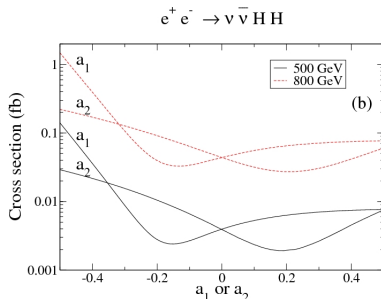
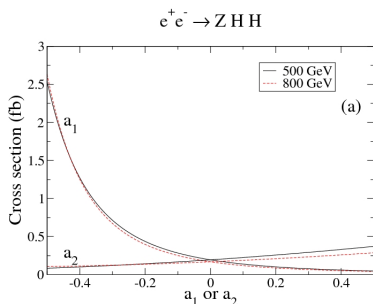
constrains by ρ : $a_1 = 4 \frac{\mu^2}{M_\Delta^4} = 0.0004 \pm_{0.0004}^{0.0007} \rightarrow$ very few deviation from SM

Any deviation observed \Rightarrow **New physics other than the triplet**

Scalar sector : ILC (3)

$$\text{Higgs pair production} / a_2 = 6\lambda_3 \frac{\mu^2}{M_\Delta^4}$$

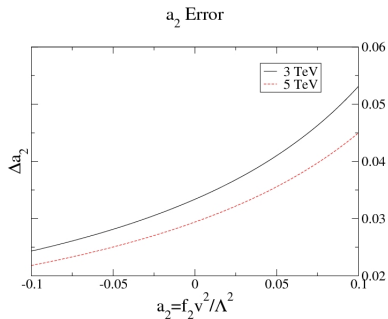
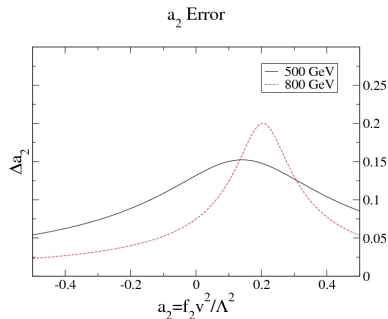
Barger & al. 03



As a_1 is strongly constrained, the only deviation is due to $a_2 = 6\lambda_3 \frac{\mu^2}{M_\Delta^4}$
→ Measurement of one parameter of the potential $\lambda_3 (\phi^\dagger\phi) (\Delta^\dagger\Delta)$

Mesurement of λ_3

Barger & al. 03



Model

$$\Delta = \begin{pmatrix} \Delta^{++} \\ \Delta^+ \\ \Delta^0 \end{pmatrix} \sim (1, 3, 2), \quad L_\Delta = -2$$

5 non-renormalizable ($d=6$) operators which effects are:

- Lepton Flavor violation
- Modification of Higgs production cross section at ILC and CLIC
- Top's physics

Seesaw Type II

- Future improvement on LFV measurement \rightarrow better knowledge on the $Y_{\Delta ij}$
- Observation of Δ^{--} at LHC \rightarrow mass scale
- Future Linear Colliders \rightarrow scalar potential, other source of NP

Seesaw Type III (Abada, Antusch, Biggio, F.B., Gavela , in progress.....)

.....Thanks for your attention.....